

Deformations and Embeddings of Compact 3D

Strictly Pseudoconvex CR Manifolds II

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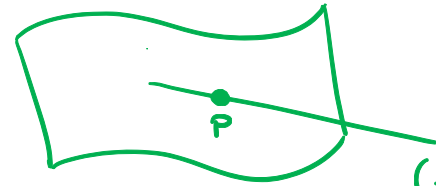
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* Joint with P. Ebenfelt

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$M^3 \subseteq \mathbb{C}^2$ real hypersurface

\leadsto CR structure (M, H, J)



(H_p, \bar{J}_p)
max. compl.
subspace

- (generically) H contact distribution

$\hookrightarrow H = \ker \theta, \theta$ 1-form
 $\theta \wedge d\theta \neq 0$

- $J: H \rightarrow H, J^2 = -\text{Id}$.

Notation: $(M, H, J) \leftrightarrow (M, T^{1,0}) \leftrightarrow (M, \bar{\partial}_b)$

$$\mathbb{C}H = T^{1,0} \oplus T^{0,1}$$

$i \quad -i$ eigenspaces
of J

Often $T^{1,0} = \text{span}\{z_1\}$

" $\bar{\partial}_b = \bar{z}_1 = \overline{z_1}$ "



$[T^{1,0}, T^{1,0}] \subseteq T^{1,0}$ trivially.

CR Structure of Unit $S^3 \subseteq \mathbb{C}^2$

$$u = 1 - |z|^2 - |w|^2, \quad \Theta = i\partial u|_{TS^3}$$

On S^3

- $\Theta = i(zd\bar{z} + wd\bar{w})|_{TS^3}$ (real)
- $d\Theta = i(dz\wedge d\bar{z} + dw\wedge d\bar{w})$
- $Z_1 = \bar{w}\partial_z - \bar{z}\partial_w$
- $Z_{\bar{1}} = w\partial_{\bar{z}} - z\partial_{\bar{w}}$
- $T = i(z\partial_z + w\partial_w) - i(\bar{z}\partial_{\bar{z}} + \bar{w}\partial_{\bar{w}})$
(Reeb v.f. of Θ)

Abstract Deformations

- Gray's theorem: might as well fix H .
- Eliashberg '89: might as well assume H is standard on S^3 .
("tight")

\leadsto Only vary J (\Leftrightarrow keep $\text{span}\{z_1, z_{\bar{1}}\}$ fixed)

$$\hat{z}_1 = \frac{1}{\sqrt{1-|\varphi|^2}} (z_1 + \underline{\varphi_1}^T z_{\bar{1}})$$

$$\hat{\theta}_1 = \frac{1}{\sqrt{1-|\varphi|^2}} (\theta_1 - \underline{\varphi_{\bar{1}}} \theta^T)$$

conj.

$|\varphi|^2 < 1$
(pointwise)

We get a parametrization of all CR structures on H in this way (up to conjugation).

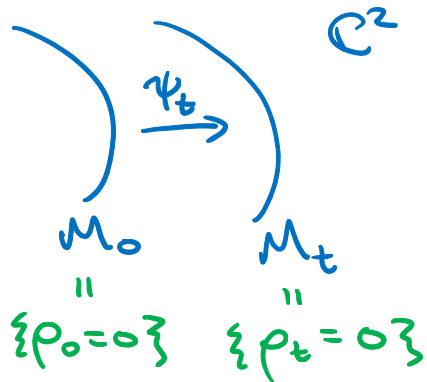
Question: When does $\varphi = \varphi_1^T$ correspond to an actual deformation of S^3 inside \mathbb{C}^2 ?

→ An embedding problem!

Remarks:

- Lempert Stability: Small deformations of S^3 (or $M^3 \subseteq \mathbb{C}^2$) that are embeddable in \mathbb{C}^N are embeddable in \mathbb{C}^2 so we always consider embeddability in \mathbb{C}^2 .
- In higher dimensions every compact (integrable) CR manifold is embeddable (Boutet de Monvel '75).
- Integrability condition:
$$\bar{\partial}_b \bar{\varphi} + \bar{\varphi} \wedge \bar{\varphi} = 0.$$

Solution at the linearized level:



Pull back by $\psi_t: M_0 \rightarrow M_t$
 (contact diffeo^m)
 to get a family of CR structures
 on M_0 .

$$\leadsto z_i^t = \frac{1}{\sqrt{1 - |\varphi_i(t)|^2}} (z_i + \underline{\varphi_i^T(t)} z_{\bar{i}})$$

$$\leadsto \dot{\varphi}_{i\bar{i}} = \left. \frac{d}{dt} \right|_{t=0} \varphi_{i\bar{i}}(t)$$

Geometric Calculation (well known): $\dot{\varphi}_{i\bar{i}}$ is a linearized
 embeddable deformation iff

$$\dot{\varphi}_{i\bar{i}} = (\nabla_i \nabla_{\bar{i}} + i A_{i\bar{i}}) f \quad \text{for some } f. \quad (*)$$

(complex)

$$\text{Re } f = \dot{\rho}|_{M_0} \quad (\text{normal velocity})$$

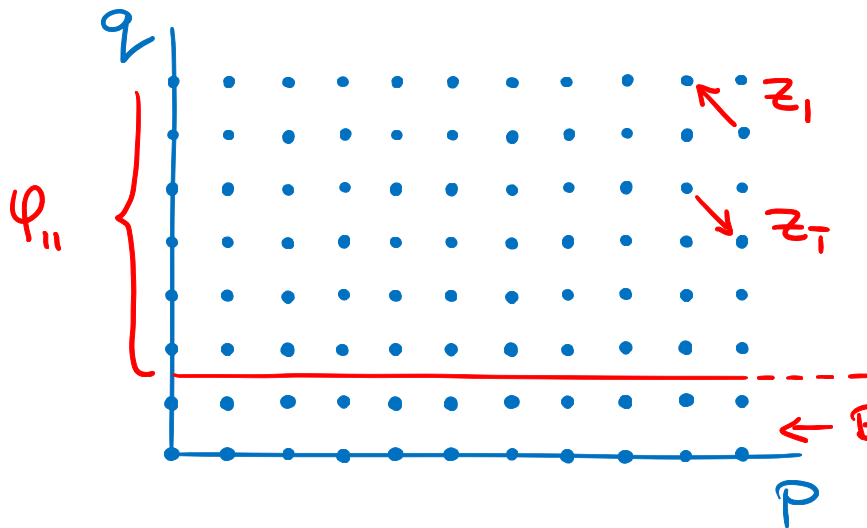
$$\left. \frac{d}{dt} \right|_{t=0} (e^{\tau_t} \rho_t)|_{M_0} = e^{\tau_0} \dot{\rho}|_{M_0}$$

On S^3 with (z_1, z_T, T) standard:

$$A_{11} = 0, \quad \omega_1' = 0$$

$$\leadsto \dot{\varphi} = (\nabla_1 \nabla_1 + i A_{11}) f = \underline{\underline{z_1 z_1}} f.$$

We understand this using spherical harmonics: $H_{p,q} \subseteq P_{p,q}$
 (e.g., $(z-w)^p (\bar{z}+\bar{w})^q \in H_{p,q}$)



$$z_1 = \bar{w} \partial_z - \bar{z} \partial_w$$

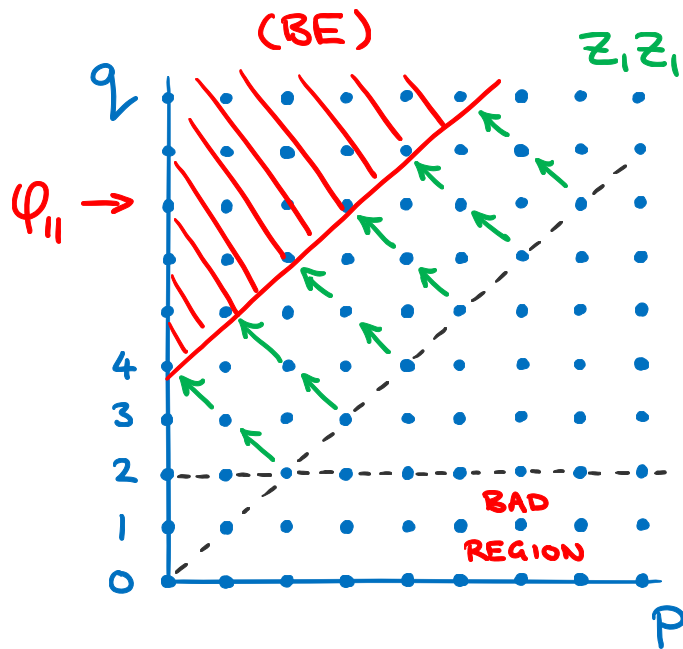
← BAD REGION

↳ Nonembeddability
is generic!

Interesting/Important Result:

Burns-Epstein, JAMS '90:

φ_{11} is embeddable if $\varphi_{11} \in \overline{\bigoplus_{q \geq p+4} H_{p,q}}$.



Good multiplicative
properties

(Using our approach we get
analyticity of the embedding
in t for $t\varphi_{11}$.)

Why 4? $L_T z_1 = -2iz_1$
 $z_1 + \underline{\varphi_{11}^T} z_{\bar{1}}$

Normalizing the CR Structure Using Contact Diffeo^ms

Q: How do contact diffeo^ms act on deformations?

Linearized level: contact (Hamiltonian) vector fields

$$X_g = gT + i(\nabla'g)Z_1 - i(\nabla\bar{g})Z_2$$

correspond to $f = ig$ ($\text{Re } f = \dot{p}|_M = 0$).

→ Linearization of natural action

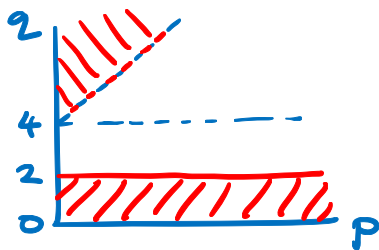
$$\{\text{Contact Diffeo's}\} \times \{\text{CR structures on } (S^3, H)\} \rightarrow \{\text{CR structures on } (S^3, H)\}$$

at $(\text{id}, J_{\text{std}})$ is

$$(g, \dot{\psi}_{11}) \mapsto \dot{\psi}_{11} + iZ_1Z_2g.$$

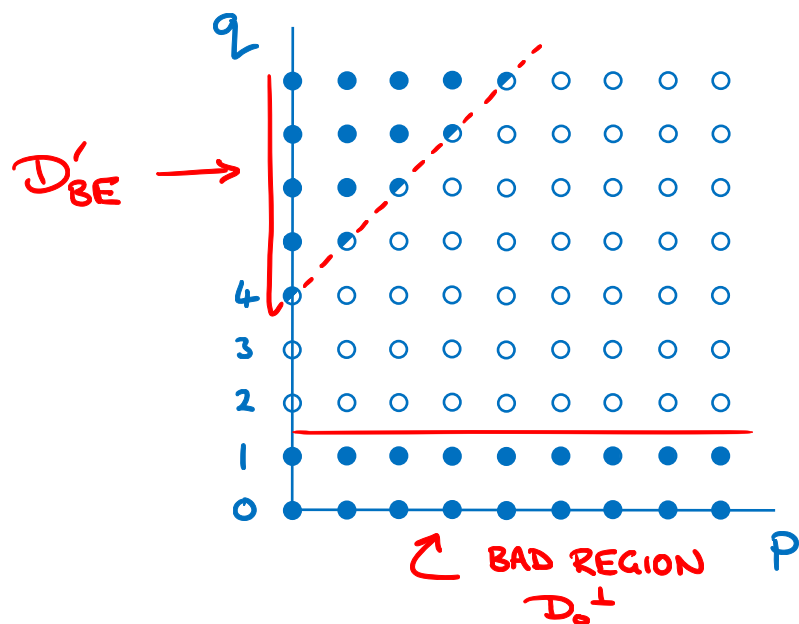
$g \in \underline{C^\infty(S^3, \mathbb{R})}$

Infinitesimal Slice:
(best choice!)



Modified Cheng-Lee Slice Theorem:

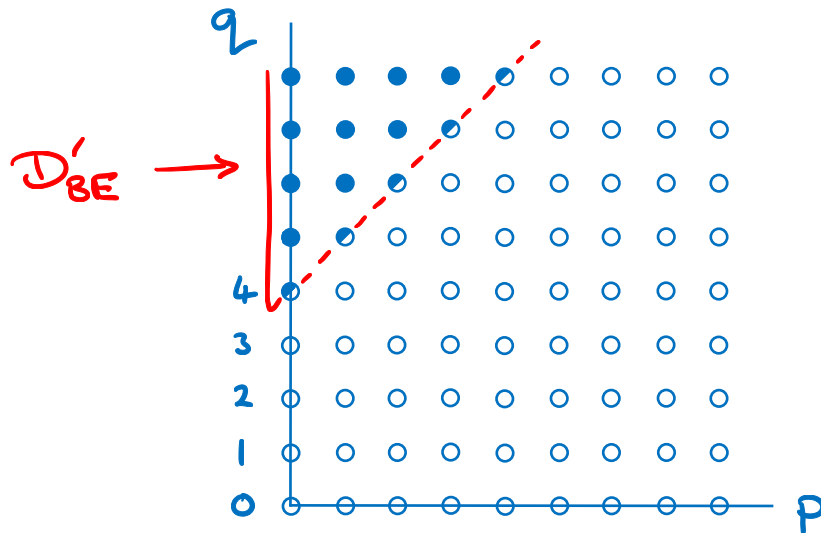
Informal Statement: Any sufficiently small deformation of the standard CR structure on (S^3, H) can be pulled back by a contact diffeo^m to one with deformation tensor φ "in"



- Proof essentially follows Cheng-Lee, AJM '95.
- Advantage is that we connect with embeddability.
- Burns-Epstein showed that the set of nonembeddable $\varphi \in (D_{BE}^\perp \oplus D_0^\perp) - D_{BE}^\perp$ is a G_δ -set.
- Epstein (Ann. Math. '98) showed this G_δ is open.
- We show it is everything (for small φ).

A Slice Theorem For Embeddable Structures:

Slice:



* Need to use "marked" str.s to deal with noncompact symmetry gp. of std CR sphere (as in Cheng-Lee).

$$\text{Aut}(\mathbb{S}^3) \curvearrowright \mathcal{D}'_{BE}$$

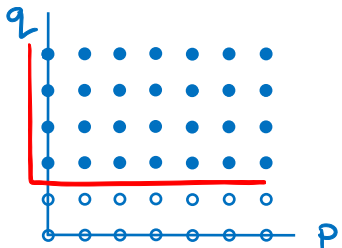
Consequence: The space of (marked) embeddable CR structures is a smooth tame Frechet submanifold near the standard CR sphere.

Relies on: Lempert (BSMF '81), Bland-Duchamp (Invent '91).
 \leadsto analytic discs, circular representation

In finite regularity: Bland (Acta '94), cf. B-D (JDA '11, JSG '14).

Q: Can we get around the need to normalize?

↗ In practice, finding the contact diffeo is a problem.

Idea: Given $\dot{\varphi}_{11} \in$  (i.e. $\varphi_{11} = z_1 \bar{z}_1 f_0$)

solve

$$(\nabla_1^t \nabla_1^t + i A_{11}(t)) f_t = \frac{\frac{d}{dt} \varphi_{11}(t)}{1 - |\varphi(t)|^2} \quad (*_t)$$

for $\varphi(t) = t\varphi + O(t^2)$ and $f_t = f_0 + O(t)$.

↖ use this to construct a family of CR embeddings realizing $\varphi(t)$.

To make this well-posed:

$$\varphi(t) = \begin{array}{c} \begin{array}{c} \uparrow \\ 2 \end{array} \\ \begin{array}{c} t\dot{\varphi} \\ + \varphi(t) \end{array} \\ \begin{array}{c} \rightarrow \\ p \end{array} \end{array}$$

$$f_t - \lambda \perp \ker z_1 \bar{z}_1 =$$

↑
const. > 0



From solutions of (\ast_t) to embeddings:

Theorem (C.-Ebenfelt, in preparation):

Given φ_t, f_t solving (\ast_t) , with $\operatorname{Re} f_t$ having strict sign,
 \exists a smooth family of embeddings $\psi_t: S^3 \rightarrow \mathbb{C}^2$ realizing φ_t
at each time t .

At $t=0$, the correct variational vector field is

$$\operatorname{Re} (f_0 \xi^0 + i \nabla' f_0 \cdot Z_1) , \quad \text{where } \xi^0 = \frac{1}{2}(\mathcal{J}T + iT).$$

But finding the variational vector field (along " $\psi_t(S^3)$ ")
in terms of f_t is much more difficult.

We refine and reverse engineer a construction in
Hirachi-Maruyama-Matsumoto (Adv. Math '17) using Fefferman's
ambient metric construction & the Graham-Lee connection.

From solutions of $(*_t)$ to embeddings:

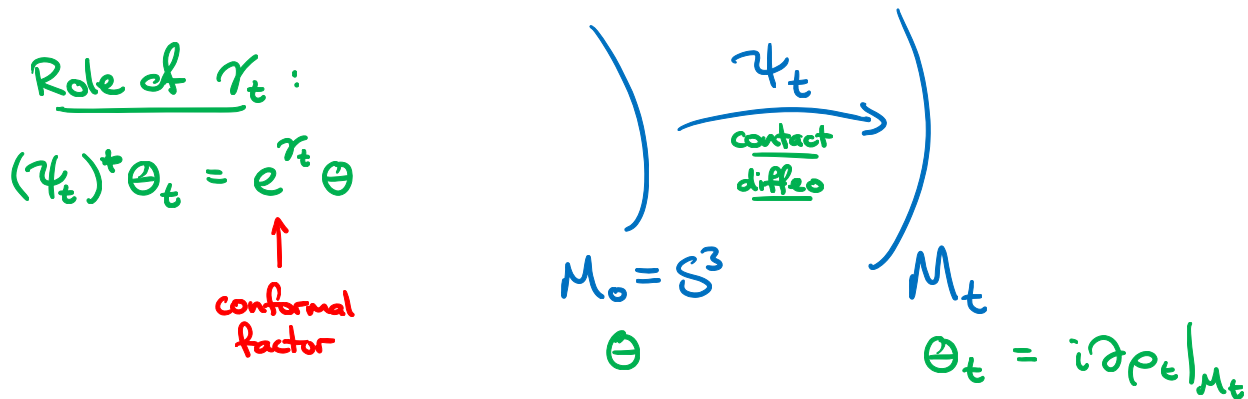
cf. Question
in Cheng '14
"Flows and a..."

Theorem (C.-Ebenfelt, in preparation):

Given φ_t, f_t solving $(*_t)$, with $\operatorname{Re} f_t$ having strict sign,
 \exists a smooth family of embeddings $\psi_t: S^3 \rightarrow \mathbb{C}^2$ realizing φ_t
 at each time t .

Transport equation:

$$\begin{cases} \dot{\psi}_t = \frac{1}{2} J \psi_t * \operatorname{Re}(e^{-2\gamma_t} f_t T) + \dots \\ \dot{\gamma}_t = -e^{-2\gamma_t} \operatorname{Re}(f_t + \dots) \end{cases}$$



What about solving $(*_t)$?

Theorem (C.-Ebenfelt, in preparation):

Given $\dot{\varphi} \in \text{image}(Z_1, Z_1)$, sufficiently small,
 $\exists \varphi(t), f_t$ solving $(*_t)$ such that

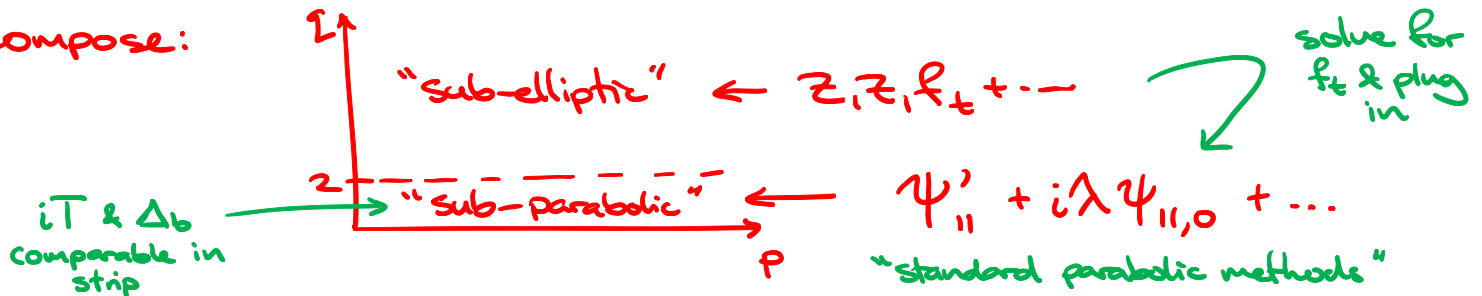
- (i) $\varphi(t) = t\dot{\varphi} + \psi(t)$ with $\psi(t) = O(t^2)$
 and $\psi(t) \in \ker(Z_1, Z_1)$,
- (ii) $\text{Re } f_t > 0$.

Indication of proof:

$$(*_t) \iff (Z_1 + \varphi_1^{-1} Z_1)^2 f_t - \varphi_1^{-1} Z_1 (Z_1 + \varphi_1^{-1} Z_1) f_t - i\lambda \varphi_{11,0} f_t = \varphi_1'$$

decompose:

(Nash-Moser)



Thank You!

